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Linearizing the Method of Conjugate Gradients
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# LINEARIZING THE METHOD OF CONJUGATE GRADIENTS* 

SERGE GRATTON ${ }^{\dagger}$, DAVID TITLEY-PELOQUIN ${ }^{\ddagger}$, PHILIPPE TOINT ${ }^{\S}$, AND JEAN TSHIMANGA ILUNGA


#### Abstract

The method of conjugate gradients (CG) is widely used for the iterative solution of large sparse systems of equations $A x=b$, where $A \in \Re^{n \times n}$ is symmetric positive definite. Let $x_{k}$ denote the $k$-th iterate of CG. In this paper we obtain an expression for $J_{k}$, the Jacobian matrix of $x_{k}$ with respect to $b$. We use this expression to obtain computable bounds on the spectral norm condition number of $x_{k}$, and to design algorithms to compute or estimate $J_{k} v$ and $J_{k}^{T} v$ for a given vector $v$. We also discuss several applications in which these ideas may be used. Numerical experiments are performed to illustrate the theory.


Key words. Conjugate Gradients Algorithm, Lanczos Algorithm, Perturbation Analysis, Linearization, Automatic Differentiation

AMS subject classifications. 65F10, 65F20, 65F22, 65F35

1. Introduction. The method of conjugate gradients (CG) of Hestenes and Stiefel [12] is widely used for the iterative solution of large sparse systems of equations $A x=b$, where $A \in \Re^{n \times n}$ is symmetric positive definite. Let $x_{k}$ denote the $k$-th iterate of CG. This iterate is a nonlinear differentiable function of $b$. In this paper we obtain an expression for the Jacobian of $x_{k}$ with respect to the right-hand side vector $b$,

$$
J_{k}=\frac{\partial x_{k}}{\partial b} \in \Re^{n \times n}
$$

The norm of the Jacobian (for any chosen norm) is by definition the absolute condition number of $x_{k}$ with respect to perturbations in $b$. We use our expression for the Jacobian to obtain computable bounds on the spectral norm condition number of $x_{k}$. We also discuss methods to compute or estimate the quantities $J_{k} v$ and $J_{k}^{T} v$ for a given vector $v$.

Sensitivity analysis of computational problems is considered by many to be a very important topic. Reference books in numerical linear algebra (such as [2], [6], [13]) typically present a sensitivity analysis of the problems they consider. For example, for the exact solution $x=A^{-1} b$, the Jacobian is

$$
\frac{\partial\left(A^{-1} b\right)}{\partial b}=A^{-1}
$$

and the condition number of $x$ in the spectral norm is $\left\|A^{-1}\right\|_{2}$, the reciprocal of the smallest singular value of $A$. Hence we are warned that the exact solution $x=A^{-1} b$

[^0]may be very sensitive to perturbations in $b$ (such as noise in the data or rounding errors in finite precision arithmetic) when $A$ has some very small singular values.

In many practical applications, however, one is not concerned with the exact solution $x=A^{-1} b$. Rather, one performs $k \ll n$ iterations of CG and accepts the resulting iterate as the computed "solution" even though it may be very far from $A^{-1} b$. For instance, when $n$ is very large, this may be the only feasible approach. Performing $k \ll n$ steps of CG is also considered a form of regularization when $A$ is very ill-conditioned; see for example the survey in [11]. Then one may wonder at quantifying this effect as a function of $b$. This is the question we consider in the present manuscript.

There has been some work done on the sensitivity of Krylov subspace methods. Kuznetsov et. al. [1, 16] obtain expressions for the condition number of a Krylov subspace $\mathcal{K}_{k}(A, b)$ with respect to perturbations in $A$ and $b$. Here, however, we are interested in the sensitivity not of a whole subspace but of only one vector in the space, namely, $x_{k}$. We also mention the papers of Greenbaum [8] and Strakoš [26] (see also [10]) who consider the sensitivity of CG iterates to changes in the eigenvalue distribution of $A$. A summary and more thorough bibliographies can be found in [19, 20]. One important aim of such work is to understand how rounding errors in finite precision arithmetic affect the convergence of the algorithm. Here our motivation is different: we are interested in applications in which $b$ is an observation vector greatly contaminated by noise, and in which very few iterations of CG are performed. In such cases, a sensitivity analysis of $x_{k}$ with respect to perturbations in $b$ is certainly relevant.

One application of this work is in statistics. Consider the linear model $K x=c$, where $K$ has full column rank. It is well known (see, e.g., [14, §17]) that given an observation $\tilde{c}=c+v$ with $v \sim(0, \Sigma), \Sigma$ non-singular, the least-squares estimator

$$
x_{L S}=\arg \min _{x}\|\tilde{c}-K x\|_{\Sigma^{-1}}=\left(K^{T} \Sigma^{-1} K\right)^{-1} K^{T} \Sigma^{-1} \tilde{c}
$$

is the best linear unbiased estimator of $x$, and

$$
\operatorname{cov}\left\{x_{L S}-x\right\}=\left(K^{T} \Sigma^{-1} K\right)^{-1}
$$

But if $x_{k}$ is the $k$-th iterate of CG applied to the normal equations of the above leastsquares problem, what is the covariance of $x_{k}-x$, and how does it compare to that of $x_{L S}-x$ ? This is sometimes called the partial least-squares problem in the statistics literature; see for example [3] and the references therein. If we know the Jacobian of $x_{k}$, we have to first order

$$
\operatorname{cov}\left\{x_{k}-x\right\} \approx J_{k} \Sigma J_{k}^{T}
$$

See e.g. [5,21] and the references therein for some data assimilation applications in which knowledge of $\operatorname{cov}\left\{x_{k}-x\right\}$ would be useful.

Another potential application is to solving linear systems with multiple righthand sides $A x=b_{i}, i=1,2, \ldots$ Suppose that after $k$ steps of CG applied to $A x=b_{i}$ we have computed $x_{k}\left(b_{i}\right)$ and $J_{k}\left(b_{i}\right)$. To first order

$$
x_{k}\left(b_{i+1}\right) \approx x_{k}\left(b_{i}\right)+J_{k}\left(b_{i}\right) \cdot\left(b_{i+1}-b_{i}\right),
$$

which is an ideal starting vector for CG applied to $A x=b_{i+1}$. If $J_{k} \approx A^{-1}$, the Jacobian can be used as a preconditioner for subsequent solves with $A$ and to estimate the energy norm of the error,

$$
\left\|\epsilon_{k}\right\|_{A}^{2}=\epsilon_{k}^{T} A \epsilon_{k}=r_{k}^{T} A^{-1} r_{k} \approx r_{k}^{T} J_{k} r_{k}
$$

where $\epsilon_{k}=A^{-1} b-x_{k}$ is the error and $r_{k}=b-A x_{k}$ the residual. In each of the above applications, in practise we would more likely use an approximation to $J_{k}$ with which matrix-vector multiplications can be performed cheaply, rather than $J_{k}$ itself. We discuss a few such approximations.

The rest of this paper is organized as follows. In Section 2 we introduce the Lanczos [17] and CG algorithms. In Section 3 we obtain an expression for $J_{k}$, the Jacobian of $x_{k}$ with respect to $b$. We also give bounds on the normwise relative error between $J_{k}$ and $A^{-1}$ and on $\left\|J_{k}\right\|_{2}$. We discuss methods to compute or estimate $J_{k} v$ and $J_{k}^{T} v$ for a given vector $v$ in Section 4. In Section 5 we present numerical experiments to illustrate the theory, and we conclude with a discussion in Section 6.
2. The Lanczos and Conjugate Gradients algorithms. We start by reviewing some known facts about the Lanczos algorithm and its relation to CG. We assume exact arithmetic. The effects of rounding errors in floating point arithmetic are discussed briefly in Section 6. For a more thorough treatment of these topics, including implementation details, see for example [2], [4], [9], [19].

The Lanczos algorithm computes an orthogonal tridiagonalization of the symmetric matrix $A \in \Re^{n \times n}$ column by column, starting from an arbitrary normalized vector $v_{1}$. After $k$ steps the algorithm produces $V_{k}=\left[v_{1}, \ldots v_{k}\right] \in \Re^{n \times k}$ with orthonormal columns and

$$
T_{k}=\left[\begin{array}{cccc}
\alpha_{1} & \beta_{2} & &  \tag{2.1}\\
\beta_{2} & \alpha_{2} & \ddots & \\
& \ddots & \ddots & \beta_{k} \\
& & \beta_{k} & \alpha_{k}
\end{array}\right], \quad \widetilde{T}_{k}=\left[\begin{array}{c}
T_{k} \\
\beta_{k+1} e_{k}^{T}
\end{array}\right],
$$

such that

$$
\begin{equation*}
A V_{k}=V_{k} T_{k}+\beta_{k+1} v_{k+1} e_{k}^{T}=V_{k+1} \widetilde{T}_{k} \tag{2.2}
\end{equation*}
$$

(Here $e_{k}$ is the $k$-th standard basis vector, not to be confused with the error $\epsilon_{k}$ below.) The columns of $V_{k}$ form an orthonormal basis for the Krylov subspace

$$
\mathcal{K}_{k}\left(A, v_{1}\right)=\operatorname{span}\left\{v_{1}, A v_{1}, \ldots, A^{k-1} v_{1}\right\} .
$$

Starting from an initial guess $x_{0}$ with residual $r_{0}=b-A x_{0}$, CG produces a sequence of iterates satisfying

$$
\begin{aligned}
& x_{k} \in x_{0}+\mathcal{K}_{k}\left(A, r_{0}\right), \\
& r_{k}=b-A x_{k} \in A \mathcal{K}_{k}\left(A, r_{0}\right), \quad r_{k} \perp \mathcal{K}_{k}\left(A, r_{0}\right) .
\end{aligned}
$$

Define $\beta_{1}=\left\|r_{0}\right\|_{2}$ and start the Lanczos algorithm with $v_{1}=r_{0} / \beta_{1}$. Then the iterate $x_{k}$ and residual $r_{k}$ from CG can be written as

$$
\begin{align*}
x_{k} & =x_{0}+V_{k} T_{k}^{-1} \beta_{1} e_{1}  \tag{2.3}\\
r_{k} & =r_{0}-A V_{k} T_{k}^{-1} \beta_{1} e_{1}=-\beta_{1} \beta_{k+1} e_{k}^{T} T_{k}^{-1} e_{1} v_{k+1}
\end{align*}
$$

We will also use the fact that $\mathcal{K}_{k}\left(A, r_{0}\right)=\operatorname{Range}\left(V_{k}\right)=\operatorname{Range}\left(K_{k}\right)$, where

$$
\begin{equation*}
K_{k}=\left[r_{0}, \ldots, A^{k-1} r_{0}\right]=\left[A \epsilon_{0}, \ldots, A^{k} \epsilon_{0}\right] \tag{2.4}
\end{equation*}
$$

and $\epsilon_{0}=A^{-1} b-x_{0}$ is the initial error.

It can also be useful to think of CG in terms of polynomials:

$$
\begin{equation*}
x_{k}=x_{0}+\zeta_{k-1}(A) r_{0} \tag{2.5}
\end{equation*}
$$

where $\zeta_{k-1}(A)$ is a polynomial of degree at most $k-1$. Then the error $\epsilon_{k}=A^{-1} b-x_{k}$ and the residual $r_{k}=b-A x_{k}$ satisfy

$$
\begin{equation*}
\epsilon_{k}=\rho_{k}(A) \epsilon_{0}, \quad r_{k}=\rho_{k}(A) r_{0}, \quad \rho_{k}(A)=I-A \zeta_{k-1}(A) \tag{2.6}
\end{equation*}
$$

Below we outline some properties of the polynomial $\rho_{k}$ to be used in later sections. Let $\Pi_{k}$ denote the set of polynomials of degree at most $k$. It is well known that

$$
\begin{equation*}
\left\|\epsilon_{k}\right\|_{A}=\left\|\rho_{k}(A) \epsilon_{0}\right\|_{A}=\min _{\substack{\rho \in \Pi_{k} \\ \rho(0)=1}}\left\|\rho(A) \epsilon_{0}\right\|_{A} \tag{2.7}
\end{equation*}
$$

Let $\mu_{k}^{(j)}, j=1, \ldots, k$, denote the eigenvalues of $T_{k}$, known as Ritz values. These are the roots of $\rho_{k}$ (see e.g. [4, §2.4], $[25, \S 2]$ ) which can therefore be written in the form

$$
\begin{equation*}
\rho_{k}(\lambda)=\prod_{j=1}^{k}\left(1-\frac{\lambda}{\mu_{k}^{(j)}}\right) . \tag{2.8}
\end{equation*}
$$

In the following lemma we give another characterization of $\rho_{k}$. Similar ideas are used in $[7, \S 1]$ and $[18, \S 6]$.

Lemma 2.1. Let the spectral decomposition of $A$ be

$$
A=Q \Lambda Q^{T}, \quad Q^{-1}=Q^{T}, \quad \Lambda=\operatorname{diag}\left(\lambda_{i}\right), \quad 0<\lambda_{1} \leq \cdots \leq \lambda_{n}
$$

and define

$$
L_{k}=\left[\begin{array}{ccc}
\lambda_{1} & \ldots & \lambda_{1}^{k}  \tag{2.9}\\
\vdots & & \vdots \\
\lambda_{n} & \ldots & \lambda_{n}^{k}
\end{array}\right], \quad w=\Lambda^{1 / 2} Q^{T} \epsilon_{0}, \quad W=\operatorname{diag}\left(w_{i}\right), \quad e=\left[\begin{array}{c}
1 \\
\vdots \\
1
\end{array}\right] .
$$

Write the polynomial $\rho_{k}$ in (2.6) as

$$
\begin{equation*}
\rho_{k}(\lambda)=1+\sum_{i=1}^{k} \tau_{i} \lambda^{i}, \quad t_{k}=\left[\tau_{1}, \ldots, \tau_{k}\right]^{T} \tag{2.10}
\end{equation*}
$$

Then provided $C G$ has not converged to the exact solution $A^{-1} b$ before step $k$, the matrix $L_{k}^{T} W^{2} L_{k}$ is non-singular and the vector of coefficients $t_{k}$ satisfies

$$
\begin{equation*}
t_{k}=-\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} W^{2} e \tag{2.11}
\end{equation*}
$$

Proof. If the matrix $L_{k}^{T} W^{2} L_{k}$ is singular there exists a nonzero vector $v \in \Re^{k}$ such that $W L_{k} v=0$. This implies, for all indices $i \in[1, k]$, either $w_{i}=0$ or

$$
v_{1} \lambda_{i}+v_{2} \lambda_{i}^{2}+\cdots+v_{k} \lambda_{i}^{k}=0
$$

Let $j$ denote the smallest index such that $v_{j} \neq 0$, and without loss of generality scale $v$ such that $v_{j}=1$. Then either $w_{i}=0$ or

$$
1+v_{j+1} \lambda_{i}+\cdots+v_{k} \lambda_{i}^{k-j}=0
$$

In other words, there is a polynomial $\tilde{\rho}_{k-j}$ of degree at most $k-j$ such that $\tilde{\rho}_{k-j}(0)=1$ and $\tilde{\rho}_{k-j}\left(\lambda_{i}\right)=0$ for all $i$ satisfying $w_{i} \neq 0$. Note from the definition of $w$ in (2.9) that $w_{i}=0$ implies $e_{i}^{T} Q \epsilon_{0}=0$. Thus, we have

$$
\tilde{\rho}_{k-j}(A) \epsilon_{0}=Q \tilde{\rho}_{k-j}(\Lambda) Q^{T} \epsilon_{0}=0
$$

The above and (2.7) implies that $\epsilon_{k-j}=0$, i.e., $x_{k-j}=A^{-1} b$. Therefore, so long as CG has not converged to the exact solution, the matrix $L_{k}^{T} W^{2} L_{k}$ cannot be singular.

Now from (2.7),

$$
\left\|\rho_{k}(A) \epsilon_{0}\right\|_{A}=\min _{\substack{\rho \in \Pi_{k} \\ \rho(0)=1}}\|\rho(\Lambda) w\|_{2}^{2}=\min _{\substack{\rho \in \Pi_{k} \\ \rho(0)=1}}\|W \rho(\Lambda) e\|_{2}^{2}=\min _{t}\left\|W\left(e+L_{k} t\right)\right\|_{2}^{2}
$$

In other words, the vector of coefficients $t_{k}$ in (2.10) is the solution of a weighted linear least-squares problem:

$$
t_{k}=\arg \min _{t}\left\|W\left(e+L_{k} t\right)\right\|_{2}^{2}=-\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} W^{2} e
$$

Lemma 2.1 gives a convenient expression for the coefficients of the polynomial $\rho_{k}$ in (2.6) in terms of the eigenvalues and eigenvectors of $A$ and the initial error $\epsilon_{0}$. We use this expression to obtain our main result in the next section.
3. The Jacobian of $x_{k}$. Let $x_{k}$ denote the $k$-th iterate of CG. The following theorem gives an expression for $J_{k}$, the Jacobian of $x_{k}$ with respect to the right-hand side vector $b$, in terms of the matrices $V_{k}$ and $T_{k}$ and the polynomial $\rho_{k}$ defined in Section 2. Throughout we assume that CG has not converged to the exact solution $A^{-1} b$ at step $k$, otherwise trivially $J_{k}=A^{-1}$.

Theorem 3.1. Let $x_{k}$ be the $k$-th iterate of $C G$ applied to $A x=b$ starting from $x_{0}$. In the notation of Section 2,

$$
\begin{equation*}
J_{k}=A^{-1}\left[I-\rho_{k}(A)\right]+2 V_{k} T_{k}^{-1} V_{k}^{T} \rho_{k}(A) \tag{3.1}
\end{equation*}
$$

Proof. Because $x_{k}=A^{-1} b-\epsilon_{k}$, we have $\partial x_{k}=A^{-1} \partial b-\partial \epsilon_{k}$, where

$$
\begin{aligned}
& \epsilon_{k}=\rho_{k}(A) \epsilon_{0}=\left(I+\sum_{i=1}^{k} \tau_{i} A^{i}\right) \epsilon_{0}, \\
& \partial \epsilon_{k}=\left(\sum_{i=1}^{k} \partial \tau_{i} A^{i}\right) \epsilon_{0}+\rho_{k}(A) \partial \epsilon_{0}=K_{k} \partial t_{k}+\rho_{k}(A) A^{-1} \partial b,
\end{aligned}
$$

and $K_{k}$ is defined in (2.4). Thus

$$
\begin{equation*}
\partial x_{k}=A^{-1}\left[I-\rho_{k}(A)\right] \partial b-K_{k} \partial t_{k} \tag{3.2}
\end{equation*}
$$

We now obtain an expression for the differential $\partial t_{k}$. Recall from Lemma 2.1 that

$$
\begin{equation*}
t_{k}=-\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} W^{2} e \tag{3.3}
\end{equation*}
$$

where only $W$ in the right-hand side depends on $b$. Therefore, the differential of $t_{k}$ is

$$
\begin{aligned}
\partial t_{k} & =\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} 2 W(\partial W) L_{k}\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} W^{2} e-\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} 2 W(\partial W) e \\
& =-2\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} W(\partial W)\left(L_{k} t_{k}+e\right) \\
& =-2\left(L_{k}^{T} W^{2} L_{k}\right)^{-1} L_{k}^{T} W \operatorname{diag}\left(\left[L_{k} t_{k}+e\right]_{i}\right) \partial w \\
& =-2\left(W L_{k}\right)^{\dagger} \rho_{k}(\Lambda) \Lambda^{-1 / 2} Q^{T} \partial b
\end{aligned}
$$

where we have used the facts that $\partial\left(M^{-1}\right)=-M^{-1}(\partial M) M^{-1}$ for any nonsingular matrix $M$ and

$$
\partial w=\Lambda^{1 / 2} Q^{T} \partial \epsilon_{0}=\Lambda^{-1 / 2} Q^{T} \partial b
$$

which follows from the definition (2.9). Now because

$$
\begin{equation*}
K_{k}=\left[A \epsilon_{0}, \ldots, A^{k} \epsilon_{0}\right]=Q \Lambda^{-1 / 2}\left[\Lambda w, \Lambda^{2} w, \ldots \Lambda^{k} w\right]=Q \Lambda^{-1 / 2} W L_{k} \tag{3.4}
\end{equation*}
$$

we have

$$
\begin{equation*}
K_{k} \partial t_{k}=-2 Q \Lambda^{-1 / 2}\left(W L_{k}\right)\left(W L_{k}\right)^{\dagger} \rho_{k}(\Lambda) \Lambda^{-1 / 2} Q^{T} \partial b \tag{3.5}
\end{equation*}
$$

Using (3.4), Range $\left(V_{k}\right)=\operatorname{Range}\left(K_{k}\right)$, and $V_{k}^{T} A V_{k}=T_{k}$, we obtain

$$
\left(W L_{k}\right)\left(W L_{k}\right)^{\dagger}=\Lambda^{1 / 2} Q^{T} K_{k}\left(K_{k}^{T} A K_{k}\right)^{-1} K_{k}^{T} Q \Lambda^{1 / 2}=\Lambda^{1 / 2} Q^{T} V_{k} T_{k}^{-1} V_{k}^{T} Q \Lambda^{1 / 2}
$$

Equation (3.1) then follows by combining the above and (3.5) into (3.2).
The formula for $J_{k}$ given in Theorem (3.1) involves the CG polynomial $\rho_{k}(A)$. Using (2.6) and (2.2) we can easily obtain equivalent expressions involving the polynomial $\zeta_{k-1}(A)$ :

$$
\begin{align*}
J_{k} & =\zeta_{k-1}(A)+2 V_{k} T_{k}^{-1} V_{k}^{T} \rho_{k}(A)  \tag{3.6a}\\
& =2 V_{k} T_{k}^{-1} V_{k}^{T}+\left(I-2 V_{k} T_{k}^{-1} V_{k}^{T} A\right) \zeta_{k-1}(A)  \tag{3.6b}\\
& =2 V_{k} T_{k}^{-1} V_{k}^{T}+\left(I-2 V_{k} V_{k}^{T}-2 \beta_{k+1} V_{k} T_{k}^{-1} e_{k} v_{k+1}^{T}\right) \zeta_{k-1}(A) \tag{3.6c}
\end{align*}
$$

Another interesting remark is that $J_{k}$ is generally not symmetric, since the product of the two symmetric matrices $V_{k} T_{k}^{-1} V_{k}^{T}$ and $\rho_{k}(A)$ in (3.1) is generally not symmetric.

If CG has converged to the exact solution at step $k$, i.e., if $x_{k}=A^{-1} b$, then $J_{k}=A^{-1}$. The following corollary gives an upper bound on the normwise relative error between $J_{k}$ and $A^{-1}$.

Corollary 3.2. In the notation of Theorem 3.1,

$$
\begin{equation*}
\frac{\left\|J_{k}-A^{-1}\right\|_{2}}{\left\|A^{-1}\right\|_{2}} \leq 3\left\|\rho_{k}(A)\right\|_{2} \tag{3.7}
\end{equation*}
$$

Furthermore, if $\left\|\rho_{k}(A)\right\|_{2}<1$,

$$
\begin{equation*}
\frac{\left|r_{k}^{T} J_{k} r_{k}\right|}{1+\left\|\rho_{k}(A)\right\|_{2}} \leq\left\|\epsilon_{k}\right\|_{A}^{2} \leq \frac{\left|r_{k}^{T} J_{k} r_{k}\right|}{1-\left\|\rho_{k}(A)\right\|_{2}} \tag{3.8}
\end{equation*}
$$

Proof. From (3.1) we have

$$
J_{k}-A^{-1}=\left(-A^{-1}+2 V_{k} T_{k}^{-1} V_{k}^{T}\right) \rho_{k}(A)
$$

The relationship (3.7) follows by takings norms and using the fact that $\left\|V_{k} T_{k}^{-1} V_{k}^{T}\right\|_{2}=$ $\left\|T_{k}^{-1}\right\|_{2} \leq\left\|A^{-1}\right\|_{2}$. Recall from Section 2 that $r_{k}^{T} V_{k}=0$. Therefore,

$$
\left\|\epsilon_{k}\right\|_{A}^{2}=r_{k}^{T} A^{-1} r_{k}=r_{k}^{T}\left(J_{k}+A^{-1} \rho_{k}(A)\right) r_{k}=r_{k}^{T} J_{k} r_{k}+r_{k}^{T} A^{-1 / 2} \rho_{k}(A) A^{-1 / 2} r_{k},
$$

from which the inequalities in (3.8) follow.
We can interpret Corollary 3.2 as follows. If $\left\|\rho_{k}(A)\right\|_{2} \ll 1$ then the Jacobian $J_{k}$ is close to $A^{-1}$ and the energy norm of the error is close to $\left(r_{k}^{T} J_{k} r_{k}\right)^{1 / 2}$. From the characterization of $\rho_{k}$ in (2.7),

$$
\left\|\rho_{k}(A)\right\|_{2} \geq \frac{\left\|\rho_{k}(A) \epsilon_{0}\right\|}{\left\|\epsilon_{0}\right\|}=\frac{\left\|\epsilon_{k}\right\|}{\left\|\epsilon_{0}\right\|}
$$

where $\|\cdot\|$ can denote either the $A$-norm, the 2 -norm, or the $A^{2}$-norm, i.e., the residual 2 -norm. If $\epsilon_{0}$ has a significant component along the eigenvector of $A$ corresponding to the largest in magnitude eigenvalue of $\rho_{k}(A)$, the above lower bound is a reasonable approximation. If this is the case, and if $\left\|\epsilon_{k}\right\| /\left\|\epsilon_{0}\right\| \ll 1$, then $\left\|\rho_{k}(A)\right\|_{2} \ll 1$. However, it is certainly possible to construct examples in which $\left\|\rho_{k}(A)\right\|_{2} \gg 1$. One such example is given in Section 5.

We can also use Theorem 3.1 to obtain bounds on $\left\|J_{k}\right\|_{2}$, the spectral norm condition number of $x_{k}$ with respect to perturbations in $b$.

Corollary 3.3. In the notation of Theorem 3.1, with the polynomial $\zeta_{k-1}$ defined in (2.5),

$$
\begin{equation*}
\left\|T_{k}^{-1} e_{1}\right\|_{2} \leq\left\|J_{k}\right\|_{2} \leq 2\left\|T_{k}^{-1}\right\|_{2}+\left(1+2\|A\|_{2}\left\|T_{k}^{-1}\right\|_{2}\right)\left\|\zeta_{k-1}(A)\right\|_{2} \tag{3.9}
\end{equation*}
$$

Proof. For the lower bound, from (3.6a) we have

$$
\left\|J_{k}\right\|_{2} \geq \frac{\left\|J_{k} r_{0}\right\|_{2}}{\left\|r_{0}\right\|_{2}}=\frac{\left\|\zeta_{k-1}(A) r_{0}+2 V_{k} T_{k}^{-1} V_{k}^{T} \rho_{k}(A) r_{0}\right\|_{2}}{\left\|r_{0}\right\|_{2}}
$$

Recall that $V_{k}^{T} \rho_{k}(A) r_{0}=V_{k}^{T} r_{k}=0$. Furthermore, using (2.5), (2.3), and the fact that $\beta_{1}=\left\|r_{0}\right\|_{2}$ we have

$$
\begin{equation*}
\left\|\zeta_{k-1}(A)\right\|_{2} \geq \frac{\left\|\zeta_{k-1}(A) r_{0}\right\|_{2}}{\left\|r_{0}\right\|_{2}}=\left\|T_{k}^{-1} e_{1}\right\|_{2} \tag{3.10}
\end{equation*}
$$

The upper bound is an immediate consequence of (3.6b).
If $r_{0}$ has a significant component along the eigenvector of $A$ corresponding to the largest eigenvalue of $\zeta_{k-1}(A)$, then the lower bound in (3.10) is a reasonable approximation and $\left\|\zeta_{k-1}(A)\right\|_{2} \approx\left\|T_{k}^{-1} e_{1}\right\|_{2}$. If this is the case, (3.9) shows that both the lower bound and upper bound on the spectral norm of the Jacobian depend only on terms involving $T_{k}^{-1}$ (as opposed to $A^{-1}$ ). In such cases, the spectral norm condition number of $x_{k}$ is essentially determined by the reciprocal of the smallest Ritz value. Of course, this reasoning does not hold in the worst case, and it may happen that $\left\|J_{k}\right\|_{2} \gg\left\|T_{k}^{-1}\right\|_{2}$. Nevertheless, in typical problems, the 2-norm condition number of CG iterates is often determined by the smallest Ritz value. Numerical examples are given in Section 5.
4. Computing Matrix-Vector Products. In most practical applications, due to memory limitations it is not possible to explicitly compute and store the entire matrix $J_{k} \in \Re^{n \times n}$. In this section we show how matrix-vector products (matvecs) $J_{k} v$ and $J_{k}^{T} v$ can be computed or estimated. This can be used to estimate the spectral norm and Frobenius norm condition numbers of $x_{k}$, as follows: for any $v$ of unit length we have a lower bound $\left\|J_{k}\right\|_{2} \geq\left\|J_{k} v\right\|_{2}$, while if $v \sim(0, I)$ then $\left\|J_{k} v\right\|_{2}^{2}$ is an unbiased estimator of $\left\|J_{k}\right\|_{F}^{2}$.
4.1. Formulas involving polynomials. We can use the expressions in (3.6) to compute the matvecs $J_{k} v$ and/or $J_{k}^{T} v$. There are several mathematically (but not numerically) equivalent methods of doing this. The computation of a degree- $k$ polynomial in $A$ times a vector using Horner's method involves $k-1$ matvecs with $A$. This is essentially as expensive as performing $k-1$ extra iterations of CG. Other algorithms for computing $\rho_{k}(A) v$ involve as few as $\sqrt{k}$ matvecs with $A$, but require at least $n^{2}$ storage; see, e.g., $[13, \S 5]$. This is much likely too expensive to be practical in large sparse applications. Even if we are willing to perform $k$ extra matvecs with $A$, the accurate computation of $\rho_{k}(A) v$ and/or $\zeta_{k-1}(A) v$ in floating point arithmetic is very challenging, particularly for larger values of $k$. (See below.)

Nevertheless, the formulas in (3.6) can be used as a starting point to obtain cheap estimates of $J_{k} v$ and $J_{k}^{T} v$. In the following estimates the matrices $V_{k}$ and $T_{k}$ are explicitly used. Therefore, these estimates really apply to the Lanczos algorithm for solving symmetric positive definite $A x=b$ (which is mathematically equivalent to CG but in which one explicitly computes and stores $V_{k}$ and $T_{k}$ ).

A natural idea is to estimate $J_{k}$ by replacing $A$ with an estimate of $A$ in (3.6). Using $A \approx V_{k} T_{k} V_{k}^{T}$ we obtain

$$
\begin{aligned}
J_{k} & =\zeta_{k-1}(A)+2 V_{k} T_{k}^{-1} V_{k}^{T} \rho_{k}(A) \\
& \approx \zeta_{k-1}\left(V_{k} T_{k} V_{k}^{T}\right)+2 V_{k} T_{k}^{-1} V_{k}^{T} \rho_{k}\left(V_{k} T_{k} V_{k}^{T}\right) \\
& =V_{k} \zeta_{k-1}\left(T_{k}\right) V_{k}^{T}+2 V_{k} T_{k}^{-1} \rho_{k}\left(T_{k}\right) V_{k}^{T} \\
& =V_{k} T_{k}^{-1}\left[I-\rho_{k}\left(T_{k}\right)\right] V_{k}^{T}+2 V_{k} T_{k}^{-1} \rho_{k}\left(T_{k}\right) V_{k}^{T}=V_{k} T_{k}^{-1} V_{k}^{T} .
\end{aligned}
$$

The last equality follows from the fact that $\rho_{k}\left(T_{k}\right)=0$, since the eigenvalues of $T_{k}$ are the zeros of $\rho_{k}$. If $A \approx V_{k} T_{k} V_{k}^{T}$ is a good approximation then hopefully so is $J_{k} \approx V_{k} T_{k}^{-1} V_{k}^{T}$, but this depends on the sensitivity of the polynomials $\zeta_{k-1}(A)$ and $\rho_{k}(A)$ to perturbations in $A$. See for example [15, §6.1] for one such perturbation result. Also note that if $v \perp \operatorname{Range}\left(V_{k}\right)$, clearly $V_{k} T_{k}^{-1} V_{k}^{T} v=0$ may be very far from $J_{k} v$.

Another simple idea is to replace $\rho_{k}(A)$ and/or $\zeta_{k-1}(A)$ in (3.6) with lower-degree polynomials. For example, from (3.6c) we can approximate

$$
\begin{equation*}
J_{k} v \approx f_{k, d^{*}}=2 V_{k} T_{k}^{-1} V_{k}^{T} v+\left(I-2 V_{k} V_{k}^{T}-2 \beta_{k+1} V_{k} T_{k}^{-1} e_{k} v_{k+1}^{T}\right) \zeta_{d-1}(A) v, \tag{4.1}
\end{equation*}
$$

where $d=\min \left\{k, d^{*}\right\}$ for some small index $d^{*}$. At step $k=d^{*}$, we compute the vector $\zeta_{d^{*}-1}(A) v$, store it, and reuse it for the following iterations. Computing the resulting approximation to $J_{k} v$ for all $k>d^{*}$ no longer requires any matvecs with $A$. Unfortunately, we have not found a way to exploit this idea to compute $J_{k}^{T} v$, since in this case the vector

$$
\left(I-2 V_{k} V_{k}^{T}-2 \beta_{k+1} v_{k+1} e_{k}^{T} T_{k}^{-1} V_{k}^{T}\right) v
$$

to be multiplied by $\zeta_{d-1}(A)$ changes at every iteration.

It is known that the polynomial $\zeta_{k-1}$ follows a three-term recurrence relation (see [19, Lemma 2.2.5]). For completeness we state it here in our notation. Recall that the scalars $\alpha_{k}$ and $\beta_{k}$ are the entries of $T_{k}$ in (2.1). Denote

$$
\eta_{k}=\frac{\alpha_{k}}{\beta_{k+1}} \frac{\left\|r_{k}\right\|_{2}}{\left\|r_{k-1}\right\|_{2}}
$$

Then

$$
\begin{aligned}
& \zeta_{0}(\lambda)=\alpha_{1}^{-1} \\
& \zeta_{1}(\lambda)=\eta_{2}\left(\alpha_{1}^{-1}+\alpha_{2}^{-1}-\alpha_{1}^{-1} \alpha_{2}^{-1} \lambda\right) \\
& \zeta_{k-1}(\lambda)=\alpha_{k}^{-1} \eta_{k}+\eta_{k}\left(1-\alpha_{k}^{-1} \lambda\right) \zeta_{k-2}(\lambda)+\left(1-\eta_{k}\right) \zeta_{k-3}(\lambda), \quad k=3,4, \ldots
\end{aligned}
$$

One can also derive an equivalent relation that uses the CG coefficients $\mu_{k}$ and $\nu_{k}$ (in Algorithm 1 below) instead of the Lanczos coefficients $\alpha_{k}$ and $\beta_{k}$ (see [19, $\left.\S 2.6\right]$ ) but because (3.6c) and (4.1) require $V_{k}$ and $T_{k}$ from the Lanczos algorithm we consider the recurrence as stated above. This immediately gives a three-term recurrence for computing the vectors $\zeta_{k-1}(A) v$. Each step in the recurrence requires one matvec with $A$. This can be done on the fly.

It is also possible to compute $\zeta_{k-1}(A) v$ from its coefficients in the monomial basis. We first compute the Ritz values, which are the roots of $\rho_{k}$. We can then interpolate at the points $\left\{(0,1),\left(\mu_{k}^{(1)}, 0\right), \ldots,\left(\mu_{k}^{(k)}, 0\right)\right\}$ to obtain the coefficients of $\rho_{k}$. From this and (2.6) we immediately obtain the coefficients of $\zeta_{k-1}$, from which we can compute $\zeta_{k-1}(A) v$ using Horner's method.
4.2. Automatic differentiation. For a given vector $v$, the matvecs $J_{k} v$ and $J_{k}^{T} v$ can also be computed directly using automatic differentiation techniques. (See for example [22, $\S 7.2]$ for an introduction to automatic differentiation.) In the following we outline how to do this.

```
Algorithm 1 The standard CG iterations
    Given \(A, b\), and \(x_{0}\)
    \(r_{0}=b-A x_{0}\)
    \(p_{0}=r_{0}\)
    \(k=1\)
    while stopping criterion not satisfied do
        \(\mu_{k}=r_{k-1}^{T} r_{k-1} / p_{k-1}^{T} A p_{k-1}\)
        \(x_{k}=x_{k-1}+\mu_{k} p_{k-1}\)
        \(r_{k}=r_{k-1}-\mu_{k} A p_{k-1}\)
        \(\nu_{k}=r_{k}^{T} r_{k} / r_{k-1}^{T} r_{k-1}\)
        \(p_{k}=r_{k}+\nu_{k} p_{k-1}\)
        \(k=k+1\)
    end while
    end
```

Algorithm 1 defines the standard CG iterations, presented essentially as in [19]. We can recast one full iteration of CG in terms of the action of operators. First we
replace the scalars $\mu_{k}$ and $\nu_{k}$ by their corresponding expressions

$$
\begin{aligned}
\mu_{k} & =\frac{r_{k-1}^{T} r_{k-1}}{p_{k-1}^{T} A p_{k-1}} \\
\nu_{k} & =\frac{r_{k}^{T} r_{k}}{r_{k-1}^{T} r_{k-1}}=\frac{\left(r_{k-1}+\mu_{k-1} A p_{k-1}\right)^{T}\left(r_{k-1}+\mu_{k-1} A p_{k-1}\right)}{r_{k-1}^{T} r_{k-1}} \\
& =1+2 \frac{p_{k-1}^{T} A r_{k-1}}{p_{k-1}^{T} A p_{k-1}}+\frac{\left.r_{k-1}^{T} r_{k-1}\right)\left(p_{k-1}^{T} A A p_{k-1}\right)}{\left(p_{k-1}^{T} A p_{k-1}\right)^{2}} .
\end{aligned}
$$

Define the vector $z_{k}=\left[x_{k}^{T}, r_{k}^{T}, p_{k}^{T}\right]^{T} \in \Re^{3 n}$. For the initial step (lines 2-4) we have

$$
z_{0}=F_{0}(b)=\left[\begin{array}{c}
x_{0} \\
b-A x_{0} \\
b-A x_{0}
\end{array}\right], \quad \frac{\partial z_{0}}{\partial b}=\left[\begin{array}{c}
0 \\
I_{n} \\
I_{n}
\end{array}\right] .
$$

At step $k$ of CG (lines 6-8) we have

$$
z_{k}=F\left(z_{k-1}\right)=\left[\begin{array}{l}
x+\frac{r^{T} r}{p^{T} A p} p \\
r+\frac{r^{T} r}{p^{T} A p} A p \\
-r-\frac{r^{T} r}{p^{T} A p} A p+\left(1+2 \frac{p^{T} A r}{p^{T} A p}+\frac{\left(r^{T} r\right)\left(p^{T} A A p\right)}{\left(p^{T} A p\right)^{2}}\right) p
\end{array}\right]_{k-1} .
$$

Differentiating the above, we obtain

$$
\frac{\partial z_{k}}{\partial z_{k-1}}=\left[\begin{array}{lll}
I_{n} & \frac{2}{\theta} p r^{T} & \mu I_{n}-2 \frac{\mu}{\theta} p q^{T} \\
0_{n} & I_{n}+\frac{2}{\theta} q r^{T} & \mu A-2 \frac{\mu}{\theta} q q^{T} \\
0_{n} & -I_{n}-\frac{2}{\theta} q r^{T}+\frac{2}{\theta} p q^{T}+\frac{2 \tau}{\theta^{2}} p r^{T} & G
\end{array}\right]_{k-1}
$$

with $q=A p, \mu=\frac{r^{T} r}{p^{T} q}, \theta=p^{T} q, \tau=q^{T} q$, and

$$
\begin{aligned}
G=-\mu A & +\frac{2 \mu}{\theta} q q^{T}+\left(1+\frac{2 q^{T} r}{\theta}+\frac{\mu \tau}{\theta}\right) I_{n} \\
& +\frac{2}{\theta} p r^{T} A-\frac{4 q^{T} r}{\theta^{2}} p q^{T}+\frac{2 \mu}{\theta} p q^{T} A-\frac{4 \mu \tau}{\theta^{2}} p q^{T}
\end{aligned}
$$

Applying the chain rule to successive iterations we have

$$
\begin{equation*}
J_{k}=\frac{\partial x_{k}}{\partial b}=\left[I_{n}, 0,0\right] \frac{\partial z_{k}}{\partial b}=\left[I_{n}, 0,0\right]\left(\frac{\partial z_{k}}{\partial z_{k-1}}\right) \ldots\left(\frac{\partial z_{1}}{\partial z_{0}}\right)\left(\frac{\partial z_{0}}{\partial b}\right) \tag{4.2}
\end{equation*}
$$

For any vector $v$, one can compute

$$
\begin{equation*}
J_{k} v=\left[I_{n}, 0,0\right]\left(\frac{\partial z_{k}}{\partial z_{k-1}}\right) \ldots\left(\frac{\partial z_{1}}{\partial z_{0}}\right)\left(\frac{\partial z_{0}}{\partial b}\right) v \tag{4.3}
\end{equation*}
$$

on the fly by updating $v=\left(\frac{\partial z_{k}}{\partial z_{k-1}}\right) v$ at step $k$ of CG. The main cost of this operation is one matvec with $A$ for each update step. This is known as the "forward" or "direct" mode in automatic differentiation.

Computing

$$
J_{k}^{T} v=\left(\frac{\partial z_{0}}{\partial b}\right)^{T}\left(\frac{\partial z_{1}}{\partial z_{0}}\right)^{T} \ldots\left(\frac{\partial z_{k}}{\partial z_{k-1}}\right)^{T}\left[I_{n}, 0,0\right]^{T} v
$$

cannot be done on the fly, since

$$
\left(\frac{\partial z_{i}}{\partial z_{i-1}}\right)^{T} \cdots\left(\frac{\partial z_{k}}{\partial z_{k-1}}\right)^{T}\left[I_{n}, 0,0\right]^{T} v
$$

is not know at step $i$ of CG. One has to store (or recompute) all the quantities involved at every step in CG and loop through the algorithm in reverse order a posteriori. This is known as the "reverse" or "adjoint" mode.

In our experience, if one is willing to perform $k$ extra matvecs with $A$, the most accurate method of computing matvecs with $J_{k}$ is using automatic differentiation. Numerical examples are given in the next section.
5. Numerical Experiments. We provide some numerical experiments to illustrate the theory developed in the previous section.
5.1. Convergence of $\left\|J_{k}\right\|_{2}$. First we illustrate the bounds on $\left\|J_{k}\right\|_{2}$ from Section 3 by showing how the convergence of $\left\|J_{k}\right\|_{2}$ is related to that of $\left\|T_{k}^{-1}\right\|_{2}$.

Example 5.1(a) is meant to illustrate worst-case, but highly unlikely, behaviour of $\left\|J_{k}\right\|_{2}$. In this example $A$ is a diagonal $128 \times 128$ matrix with 64 eigenvalues equally spaced in $\left[10^{-3}, 10^{-2}\right]$ and 64 eigenvalues equally spaced in $\left[10^{2}, 10^{3}\right]$. The right-hand side vector is $b=[1, \ldots 1,0, \ldots, 0]^{T}$ and the iteration is started with $x_{0}=0$. Because $r_{0}$ is orthogonal to all eigenvectors corresponding to the eigenvalues in $\left[10^{2}, 10^{3}\right]$, the Lanczos algorithm fails to compute any Ritz value in this interval. Consequently, for all $k$,

$$
\begin{align*}
& \left\|\rho_{k}(\Lambda)\right\|_{2} \geq\left|\rho_{k}\left(\lambda_{n}\right)\right|=\prod_{j=1}^{k}\left|1-\frac{\lambda_{n}}{\mu_{j}}\right| \geq \prod_{j=1}^{k}\left(\frac{10^{3}}{10^{-2}}-1\right) \approx 10^{5 k}  \tag{5.1}\\
& \left\|\zeta_{k-1}(\Lambda)\right\|_{2} \geq\left|\zeta_{k-1}\left(\lambda_{n}\right)\right|=\frac{1}{\lambda_{n}}\left|1-\rho_{k}\left(\lambda_{n}\right)\right| \approx 10^{5 k-3}
\end{align*}
$$

In other words, the condition number of $x_{k}$ grows very quickly as $k$ increases, and the iterates are soon highly sensitive to perturbations in $b$. This is illustrated in Figures 5.1 and 5.2.


Fig. 5.1. Left: Polynomials $\zeta_{4}(\lambda)$ (light) and $\rho_{5}(\lambda)$ (dark) plotted vs $\lambda$ and in particular evaluated at the eigenvalues of $A$ (stars). The vertical lines show the zeros of $\rho_{5}$, i.e., the Ritz values $\mu_{5}^{(j)}, j=1, \ldots, 5$. The horizontal lines denote $\left\|\zeta_{4}(\Lambda)\right\|_{2}$ and $\left\|\rho_{5}(\Lambda)\right\|_{2}$. Right: Convergence of the Ritz values as $k$ increases. The eigenvalues of $A$ are shown at the right. No Ritz value converges to an eigenvalue in $\left[10^{2}, 10^{3}\right]$.


FIG. 5.2. Convergence of $\left\|J_{k}\right\|_{2}$ vs $k$ in $C G$ for example 5.1(a). As predicted by (5.1), $\left\|J_{k}\right\|_{2}$ increases with a slope close to 5 on the semilog plot. The plotted tolerance is $\operatorname{tol}_{k}=\epsilon\left(\|A\|_{2}\left\|x_{k}\right\|_{2}+\right.$ $\left.\|b\|_{2}\right)$, where $\epsilon$ is the machine unit roundoff.

In the previous example the Ritz values corresponding to large eigenvalues failed to converge entirely. It is well established that this is highly atypical behaviour in CG. The next examples illustrate more typical behaviour. In each case the matrix $A$ is formed via its spectral decomposition. The matrix of eigenvectors is the $Q$ factor in the QR decomposition of a random matrix, $b=e$, and $x_{0}=0$. In example 5.1(b) the eigenvalues of $A$ are logarithmically equally spaced between $10^{-4}$ and 1 , while in example 5.1 (c) there are $n-1$ eigenvalues of $A$ linearly equally spaced between 1 and 10 with one extra eigenvalue $10^{-7}$.

In these small examples we explicitly compute and store $J_{k}$ using the automatic differentiation techniques described in Section 4. For the sake of comparison we implement CG with and without reorthogonalization (of the residual vectors). When we reorthogonalize in CG we also differentiate the reorthogonalization steps of the algorithm. To compute $\left\|T_{k}^{-1}\right\|_{2}$ we separately run the Lanczos algorithm. When we reorthogonalize in CG we also reorthogonalize in Lanczos, so that $T_{k}$ is actually upper Hessenberg.

As shown in Figure 5.3, in these examples with and without reorthogonalization the spectral norm condition number essentially converges at the same rate as $\left\|T_{k}^{-1}\right\|_{2}$, which can be much smaller than $\left\|A^{-1}\right\|_{2}$ in the early iterations. When no reorthogonalization is used, the spectral norm of the Jacobian (as computed using automatic differentiation techniques) sometimes oscillates in the later iterations. We believe that this is a numerical artefact from the computation of $J_{k}$. The convergence of $\left\|J_{k}\right\|_{2}$ is typically smoother when reorthogonalization is used.
5.2. Computation of Matrix-Vector Products. We implement (3.6c) to compute $\left\|J_{k} v\right\|_{2}$ for a random vector $v \sim \mathcal{N}(0, I)$. We compute $\zeta_{k-1}(A) v$ both via the three-term recurrence and via interpolation of $\rho_{k}$ at its roots, as described in Section 4. We also implement the related estimate $\left\|f_{k, d^{*}}\right\|_{2}$ in (4.1) with $d^{*}=5$ and $d^{*}=10$. Results are plotted in Figure 5.4.

Neither approach for computing $\zeta_{k-1}(A) v$ is reliable for large values of $k$. The accuracy of this computation seems to be problem-dependent. Formula (3.6c) gave very accurate results for test problem 5.1(c), However, for problem 5.1(b), there were severe oscillations in the computed $\left\|J_{k} v\right\|_{2}$, especially when no reorthogonalization


FIG. 5.3. Convergence of $\left\|J_{k}\right\|_{2}$ vs $k$ in $C G$ for examples 5.1(b) and 5.1(c). Top: example 5.1(b). Bottom: example 5.1(c). Left: with reorthogonalization. Right: without reorthogonalization. The plotted tolerance is $\operatorname{tol}_{k}=\epsilon\left(\|A\|_{2}\left\|x_{k}\right\|_{2}+\|b\|_{2}\right)$, where $\epsilon$ is the machine unit roundoff.
was used. Even with reorthogonalization, in problem 5.1(b) the quantity $\left\|J_{k} v\right\|_{2}$ with $\zeta_{k-1}(A) v$ computed using the three-term recurrence diverged.

In (4.1), however, the quantity $\zeta_{d-1}(A) v$ is computed only for small values of $d \leq d^{*}$. In these examples, each method for computing $\zeta_{d-1}(A) v$ gave similar results, so we only plot the estimate with $\zeta_{d-1}(A) v$ computed using the three-term recurrence. The estimate $\left\|f_{k, d^{*}}\right\|_{2}$ in (4.1) is a reasonable approximation to $\left\|J_{k} v\right\|_{2}$, even for large values of $k>d^{*}$. It tends to under-estimate $\left\|J_{k} v\right\|_{2}$ in the early iterations (but less so than $\left\|V_{k} T_{k}^{-1} V_{k}^{T} v\right\|_{2}$ ) and slightly over-estimate $\left\|J_{k} v\right\|_{2}$ in the later iterations.
6. Discussion. We have performed a first order perturbation analysis for CG iterates. Our results quantify how sensitive the iterates are to perturbations in the right-hand side vector, which is of interest in many applications. In Theorem 3.1 we obtained an expression for the Jacobian of $x_{k}$ in CG in terms of the matrices $V_{k}$ and $T_{k}$ from the Lanczos algorithm and the polynomials $\rho_{k}$ in (2.6). We used this result to bound the spectral norm condition number of $x_{k}$. We showed that the condition number in the spectral norm typically converges as $\left\|T_{k}^{-1}\right\|_{2}$, the reciprocal of the smallest Ritz value, and this quantity can be much smaller in the early iterations than $\left\|A^{-1}\right\|_{2}$. This quantifies the regularization property of CG that has been observed in practise.

We have also discussed methods to compute or estimate the matvecs $J_{k} v$ and $J_{k}^{T} v$. In our experience, automatic differentiation seems to be the most reliable way to compute $J_{k} v$ and $J_{k}^{T} v$, at the cost of $k$ extra matvecs with $A$ (as well as extra storage if $J_{k}^{T} v$ is required). So far we have not found a way to compute $J_{k} v$ or $J_{k}^{T} v$


FIG. 5.4. Convergence of $\left\|J_{k} v\right\|_{2}$ and its estimates vs $k$ for examples 5.1(b) and 5.1(c). Top: example 5.1(b). Bottom: example 5.1(c). Left: with reorthogonalization. Right: without reorthogonalization. In the legend, $A D$ stands for automatic differentiation, 3TR for three-term recurrence, and INT for interpolation.
very accurately without performing $k$ extra matvecs with $A$. We have shown that a cheap and easy way to estimate $J_{k}$ is $J_{k} \approx V_{k} T_{k}^{-1} V_{k}^{T}$. We have also proposed schemes that can improve the accuracy of this estimate using fewer than $k$ matvecs with $A$. Whether or not this can be improved upon is an open question.

In deriving Theorem 3.1 we have analyzed the Lanczos and CG algorithms assuming exact arithmetic. In finite precision the relationship (2.2) holds to machine precision; however, the columns of $V_{k}$ can quickly lose their orthogonality and even their linear independence. Paige $[23,24]$ has recently showed that there exists a matrix $H_{k}$ with $\left\|H_{k}\right\|_{2} \leq \epsilon\|A\|_{2}, \epsilon$ the unit roundoff, such that the matrix $T_{k}$ obtained in finite precision satisfies

$$
\left(\left[\begin{array}{cc}
A & \\
& T_{k}
\end{array}\right]+H_{k}\right) Q_{k}=Q_{k} T_{k}+\beta_{k+1} q_{k+1} e_{k}^{T}
$$

where the columns of $Q_{k}$ are exactly orthogonal and $q_{1}=\left[v_{1}^{T}, 0\right]^{T}$. In other words, the computed $T_{k}$ is the tridiagonal matrix produced by the exact Lanczos process applied to a small perturbation of an augmented matrix $\operatorname{diag}\left(A, T_{k}\right)$ started with the augmented vector $\left[v_{1}^{T}, 0\right]^{T}$. Paige called the above the augmented backward stability of the Lanczos algorithm. In the future we intend to use this result to analyze the true sensitivity of the Lanczos and CG algorithms implemented in floating point arithmetic. (By this we do not mean the difference between $x_{k}$ and $\tilde{x}_{k}$, the iterates computed in exact and floating point arithmetic, respectively, but rather between $\tilde{x}_{k}(b)$ and $\left.\tilde{x}_{k}(b+\delta b).\right)$

We are working on extending the ideas presented in this manuscript to minimumresidual and other polynomial-based iterative methods. For CG and other methods we can also consider other derivatives such as $\frac{\partial x_{k}}{\partial x_{0}}, \frac{\partial\left\|r_{k}\right\|_{2}}{\partial r_{0}}$, etc. It may also be possible to find sparse approximations of these derivatives. Another interesting and very challenging problem which we have not considered here is the computation of the derivative of $x_{k}$ with respect to elements of $A$.

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